

Mean-Field Dynamics of Fermionic Systems

Niels Benedikter

joint work with Vojkan Jakšić, Marcello Porta, Chiara Saffirio and Benjamin Schlein

Motivation

- Many physical systems can be modelled as fermionic many-body theories.
Many non-rigorous approximation schemes exist.
- Rigorous understanding of approximations is difficult, in particular in **non-equilibrium**.

Motivation

- Many physical systems can be modelled as fermionic many-body theories.
Many non-rigorous approximation schemes exist.
- Rigorous understanding of approximations is difficult, in particular in **non-equilibrium**.
- We establish one of the most fundamental approximations: **time-dependent Hartree-Fock theory**.
- Central requirement: **quantitative understanding of 'semi-classicality'**.
- Ideas turn out to apply very generally (**extension to mixed states, derivation of Vlasov equation, . . .**).

Outline

- 1 Fermionic Many-Body Systems
- 2 Main Result: Hartree-Fock Theory is a Valid Approximation.
- 3 ... assuming certain Semi-Classical Estimates
- 4 Proof of Hartree-Fock Theory
- 5 Mixed States, Vlasov Equation

Outline

- 1** Fermionic Many-Body Systems
- 2 Main Result: Hartree-Fock Theory is a Valid Approximation.
- 3 ... assuming certain Semi-Classical Estimates
- 4 Proof of Hartree-Fock Theory
- 5 Mixed States, Vlasov Equation

Fermionic Many-Body System

- N fermions in \mathbb{R}^3 are described by wavefunction

$$\psi \in L_a^2(\mathbb{R}^{3N}), \quad \|\psi\|_2 = 1,$$

antisymmetric w. r. t. permutations.

Fermionic Many-Body System

- N fermions in \mathbb{R}^3 are described by wavefunction

$$\psi \in L_a^2(\mathbb{R}^{3N}), \quad \|\psi\|_2 = 1,$$

antisymmetric w. r. t. permutations.

- One-particle density matrix

$$\gamma_\psi := \text{tr}_{2,\dots,N} |\psi\rangle\langle\psi|.$$

Sufficient for calculating one-particle observables.

Time Evolution

- Schrödinger equation

$$i\partial_t\psi_t = H\psi_t, \quad \psi_0 \in L_a^2(\mathbb{R}^{3N}),$$

with Hamilton operator

$$H = \sum_{j=1}^N -\Delta_{x_j} + \lambda \sum_{1 \leq i < j \leq N} V(x_i - x_j).$$

Goal

Approximate one-particle density matrix γ_{ψ_t} by time-dependent Hartree-Fock equation and estimate error as $N \rightarrow \infty$.

Physical Regime

- **mean-field scaling**: high density and weak interaction, so that the particles move in averaged potential.

Physical Regime

- **mean-field scaling**: high density and weak interaction, so that the particles move in averaged potential.
- The scale of volume is fixed by the support of initial data.

Physical Regime

- **mean-field scaling**: high density and weak interaction, so that the particles move in averaged potential.
- The scale of volume is fixed by the support of initial data.
- **Pauli exclusion principle**: occupy momenta up to $|k| \simeq N^{1/3}$

$$\left\langle \sum_{j=1}^N -\Delta_{x_j} \right\rangle \simeq \sum_{j=1}^N k_j^2 \simeq N^{5/3}$$

Physical Regime

- **mean-field scaling**: high density and weak interaction, so that the particles move in averaged potential.
- The scale of volume is fixed by the support of initial data.
- **Pauli exclusion principle**: occupy momenta up to $|k| \simeq N^{1/3}$

$$\left\langle \sum_{j=1}^N -\Delta_{x_j} \right\rangle \simeq \sum_{j=1}^N k_j^2 \simeq N^{5/3}$$

- Weak, but nontrivial, interaction:

$$i\partial_t \psi_t = \left[\sum_{j=1}^N -\Delta_{x_j} + \frac{1}{N^{1/3}} \sum_{1 \leq i < j \leq N} V(x_i - x_j) \right] \psi_t.$$

Physical Regime

- Rescale time to scale of order $N^{-1/3}$

$$iN^{1/3}\partial_t\psi_t = \left[\sum_{j=1}^N -\Delta_{x_j} + \frac{1}{N^{1/3}} \sum_{1 \leq i < j \leq N} V(x_i - x_j) \right] \psi_t.$$

Physical Regime

- Rescale time to scale of order $N^{-1/3}$

$$iN^{1/3}\partial_t\psi_t = \left[\sum_{j=1}^N -\Delta_{x_j} + \frac{1}{N^{1/3}} \sum_{1 \leq i < j \leq N} V(x_i - x_j) \right] \psi_t.$$

- Introduce semi-classical parameter $\hbar := N^{-1/3}$, multiply by \hbar^2 :

Mean-field scaling is coupled to semi-classical scaling

$$i\hbar\partial_t\psi_t = \left[\sum_{j=1}^N -\hbar^2\Delta_{x_j} + \frac{1}{N} \sum_{1 \leq i < j \leq N} V(x_i - x_j) \right] \psi_t$$

$$\hbar = N^{-1/3}, \quad N \rightarrow \infty.$$

Outline

- 1 Fermionic Many-Body Systems
- 2 Main Result: Hartree-Fock Theory is a Valid Approximation.**
- 3 ... assuming certain Semi-Classical Estimates
- 4 Proof of Hartree-Fock Theory
- 5 Mixed States, Vlasov Equation

Hartree-Fock Approximation

Restrict to Slater determinants

$$\psi(x_1, x_2, \dots) = \frac{1}{\sqrt{N!}} \det(\varphi_i(x_j))_{1 \leq i, j \leq N},$$

and **optimize** the choice of the $\varphi_i \in L^2(\mathbb{R}^3)$.

Hartree-Fock Approximation

Restrict to Slater determinants

$$\psi(x_1, x_2, \dots) = \frac{1}{\sqrt{N!}} \det(\varphi_i(x_j))_{1 \leq i, j \leq N},$$

and **optimize** the choice of the $\varphi_i \in L^2(\mathbb{R}^3)$.

- Approximate time evolution

$$\frac{1}{\sqrt{N!}} \det(\varphi_{i,t}(x_j)) \simeq e^{-iHt/\hbar} \frac{1}{\sqrt{N!}} \det(\varphi_{i,0}(x_j))$$

Hartree-Fock Approximation

Restrict to Slater determinants

$$\psi(x_1, x_2, \dots) = \frac{1}{\sqrt{N!}} \det(\varphi_i(x_j))_{1 \leq i, j \leq N},$$

and **optimize** the choice of the $\varphi_i \in L^2(\mathbb{R}^3)$.

- Approximate time evolution

$$\frac{1}{\sqrt{N!}} \det(\varphi_{i,t}(x_j)) \simeq e^{-iHt/\hbar} \frac{1}{\sqrt{N!}} \det(\varphi_{i,0}(x_j))$$

- Hartree-Fock equation

$$i\hbar\partial_t\varphi_{i,t} = -\hbar^2\Delta\varphi_{i,t} + \frac{1}{N} \sum_{j=1}^N \left(V * |\varphi_{j,t}|^2 \right) \varphi_{i,t} - \frac{1}{N} \sum_{j=1}^N \left(V * (\varphi_{i,t}\overline{\varphi_{j,t}}) \right) \varphi_{j,t}$$

$$i\hbar\partial_t\gamma_t^{\text{HF}} = \left[-\hbar^2\Delta + V * \rho_t - X_t, \gamma_t^{\text{HF}} \right]$$

Validity of HF Approximation

Is γ_t^{HF} close to γ_{ψ_t} , in mean-field scaling as $N \rightarrow \infty$?

- *Narnhofer-Sewell* '81: Convergence to Vlasov equation (= semi-classical limit of HF). Analytic V .
- *Spohn* '81: more general V .
- *Erdős-Elgart-Schlein-Yau* '04: Convergence to HF for short times, $t < t_0$. Analytic V .
- *B-Porta-Schlein* '13: more general V . **Arbitrary times t .**

other physical regimes:

Bardos-Golse-Gottlieb-Mauser '03, *Fröhlich-Knowles* '11, *Pickl-Petrat* '14, *Bach-Breteaux-Petrat-Pickl-Tzaneteas* '15.

Validity of HF Approximation

Theorem (B-Porta-Schlein '13)

Let $V \in L^1(\mathbb{R}^3)$ with $\int |\hat{V}(p)|(1+|p|)^2 dp < \infty$.

Let $\{\varphi_{j,0}\}_{j=1}^\infty$ be orthonormal in $L^2(\mathbb{R}^3)$. Let $\psi_0 = \frac{1}{\sqrt{N!}} \det(\varphi_{i,0}(x_j))$.

Assume '*semi-classical commutators*'

$$\|[\hat{X}, \gamma_{\psi_0}]\|_{\text{tr}} \leq \hbar C, \quad \|[-i\hbar\nabla, \gamma_{\psi_0}]\|_{\text{tr}} \leq \hbar C.$$

Start with the same initial data ψ_0 for Schrödinger equation and HF.

Then

$$\|\gamma_{\psi_t} - \gamma_t^{HF}\|_{\text{tr}} \leq \frac{C}{N^{5/6}} e^{ce^{c|t|}}.$$

Outline

- 1 Fermionic Many-Body Systems
- 2 Main Result: Hartree-Fock Theory is a Valid Approximation.
- 3 ... assuming certain Semi-Classical Estimates**
- 4 Proof of Hartree-Fock Theory
- 5 Mixed States, Vlasov Equation

The Semi-Classical Commutators

- e. g. ground state of trapped systems (with trap to be removed to observe evolution).
- non-interacting fermions in a box

$$\varphi_{j,0}(x) = e^{ik_j x}, \quad k_j \in 2\pi\mathbb{Z}^3.$$

The Semi-Classical Commutators

- e. g. ground state of trapped systems (with trap to be removed to observe evolution).
- non-interacting fermions in a box

$$\varphi_{j,0}(x) = e^{ik_j x}, \quad k_j \in 2\pi\mathbb{Z}^3.$$

One-particle density matrix:

$$\begin{aligned} \gamma_{\psi_0}(x; y) &= \frac{1}{N} \sum_{j=1}^N \varphi_{j,0}(x) \overline{\varphi_{j,0}(y)} = \frac{1}{N} \sum_{|k| \leq cN^{1/3}} e^{ik(x-y)} \\ &\simeq \int_{|q| \leq 1} e^{iq(x-y)/\hbar} dq = \varphi\left(\frac{x-y}{\hbar}\right). \end{aligned}$$

The Semi-Classical Commutators

- e. g. ground state of trapped systems (with trap to be removed to observe evolution).
- non-interacting fermions in a box

$$\varphi_{j,0}(x) = e^{ik_j x}, \quad k_j \in 2\pi\mathbb{Z}^3.$$

One-particle density matrix:

$$\begin{aligned} \gamma_{\psi_0}(x; y) &= \frac{1}{N} \sum_{j=1}^N \varphi_{j,0}(x) \overline{\varphi_{j,0}(y)} = \frac{1}{N} \sum_{|k| \leq cN^{1/3}} e^{ik(x-y)} \\ &\simeq \int_{|q| \leq 1} e^{iq(x-y)/\hbar} dq = \varphi\left(\frac{x-y}{\hbar}\right). \end{aligned}$$

$$[\hat{X}, \gamma_{\psi_0}](x; y) = (x-y) \varphi\left(\frac{x-y}{\hbar}\right) \sim \hbar.$$

The Semi-Classical Commutators

- e. g. ground state of trapped systems (with trap to be removed to observe evolution).
- non-interacting fermions in a box

$$\varphi_{j,0}(x) = e^{ik_j x}, \quad k_j \in 2\pi\mathbb{Z}^3.$$

One-particle density matrix:

$$\begin{aligned} \gamma_{\psi_0}(x; y) &= \frac{1}{N} \sum_{j=1}^N \varphi_{j,0}(x) \overline{\varphi_{j,0}(y)} = \frac{1}{N} \sum_{|k| \leq cN^{1/3}} e^{ik(x-y)} \\ &\simeq \int_{|q| \leq 1} e^{iq(x-y)/\hbar} dq = \varphi\left(\frac{x-y}{\hbar}\right). \end{aligned}$$

$$[\hat{X}, \gamma_{\psi_0}](x; y) = (x-y) \varphi\left(\frac{x-y}{\hbar}\right) \sim \hbar.$$

- more rigorously by coherent states quantization

Outline

- 1 Fermionic Many-Body Systems
- 2 Main Result: Hartree-Fock Theory is a Valid Approximation.
- 3 ... assuming certain Semi-Classical Estimates
- 4 Proof of Hartree-Fock Theory**
- 5 Mixed States, Vlasov Equation

Particle-Hole Transformation

- embed in Fock space using 2nd quantization (for convenience)

Redefine 'particle' as 'excitation w. r. t. Slater determinant', using a unitarily implemented Bogoliubov transformation \mathbb{U}_t .

Particle-Hole Transformation

- embed in Fock space using 2nd quantization (for convenience)

Redefine 'particle' as 'excitation w. r. t. Slater determinant', using a unitarily implemented Bogoliubov transformation \mathbb{U}_t .

- new vacuum = Fermi sea

$$\left\{ 0, \dots, 0, \frac{1}{\sqrt{N!}} \det(\varphi_{i,t}(x_j)), 0, \dots \right\} \in \mathcal{F}.$$

Particle-Hole Transformation

- embed in Fock space using 2nd quantization (for convenience)

Redefine 'particle' as 'excitation w. r. t. Slater determinant', using a unitarily implemented Bogoliubov transformation \mathbb{U}_t .

- new vacuum = Fermi sea

$$\left\{ 0, \dots, 0, \frac{1}{\sqrt{N!}} \det(\varphi_{i,t}(x_j)), 0, \dots \right\} \in \mathcal{F}.$$

- $v_t =$ projection on $\text{span}\{\varphi_{1,t}, \dots, \varphi_{N,t}\}$
 $u_t = v_t^\perp$

$$v_{t,x}(y) = v_t(y; x)$$

$$\mathbb{U}_t a_x^* \mathbb{U}_t^* = a(v_{t,x}) + a^*(u_{t,x}).$$

Particle-Hole Transformation

- embed in Fock space using 2nd quantization (for convenience)

Redefine 'particle' as 'excitation w. r. t. Slater determinant', using a unitarily implemented Bogoliubov transformation \mathbb{U}_t .

- new vacuum = Fermi sea

$$\left\{ 0, \dots, 0, \frac{1}{\sqrt{N!}} \det(\varphi_{i,t}(x_j)), 0, \dots \right\} \in \mathcal{F}.$$

- $v_t =$ projection on $\text{span}\{\varphi_{1,t}, \dots, \varphi_{N,t}\}$ $v_{t,x}(y) = v_t(y; x)$
 $u_t = v_t^\perp$

$$\mathbb{U}_t a_x^* \mathbb{U}_t^* = a(v_{t,x}) + a^*(u_{t,x}).$$

- Number of excitations w. r. t. HF-evolved Slater determinant

$$\mathcal{N}_t^{\text{exc}} = \mathbb{U}_t \mathcal{N} \mathbb{U}_t^*.$$

■ Error \leq Number of Excitations

$$\|\gamma_{\psi_t} - \gamma_t^{\text{HF}}\|_{\text{tr}} \leq \frac{C}{N^{1/2}} \underbrace{\langle e^{-i\mathcal{H}t/\hbar}\psi_0, \mathcal{N}_t^{\text{exc}} e^{-i\mathcal{H}t/\hbar}\psi_0 \rangle}_{=: r(t)}.$$

- Error \leq Number of Excitations

$$\|\gamma_{\psi_t} - \gamma_t^{\text{HF}}\|_{\text{tr}} \leq \frac{C}{N^{1/2}} \underbrace{\langle e^{-i\mathcal{H}t/\hbar}\psi_0, \mathcal{N}_t^{\text{exc}} e^{-i\mathcal{H}t/\hbar}\psi_0 \rangle}_{=: r(t)}.$$

- sufficient to prove (Grönwall)

$$r'(t) \leq C_t r(t) \quad \text{with } C_t = \mathcal{O}(N^0).$$

- most terms in $r'(t)$ cancel, remaining

$$\frac{1}{\hbar N} \int dx dy V(x-y) a^*(u_{t,y}) a(u_{t,y}) a(v_{t,x}) a(u_{t,x})$$

- Error \leq Number of Excitations

$$\|\gamma_{\psi_t} - \gamma_t^{\text{HF}}\|_{\text{tr}} \leq \frac{C}{N^{1/2}} \underbrace{\langle e^{-i\mathcal{H}t/\hbar} \psi_0, \mathcal{N}_t^{\text{exc}} e^{-i\mathcal{H}t/\hbar} \psi_0 \rangle}_{=: r(t)}.$$

- sufficient to prove (Grönwall)

$$r'(t) \leq C_t r(t) \quad \text{with } C_t = \mathcal{O}(N^0).$$

- most terms in $r'(t)$ cancel, remaining

$$\frac{1}{\hbar N} \int dx dy V(x-y) a^*(u_{t,y}) a(u_{t,y}) a(v_{t,x}) a(u_{t,x})$$

- Semi-classical commutators:

$$V(x-y) = \int \hat{V}(p) e^{ip \cdot x} e^{-ip \cdot y} dp$$

$$\int dx v_{t,x} e^{ip \cdot \hat{x}} u_{t,x} = \int dx v_{t,x} [e^{ip \cdot \hat{x}}, u_t](\cdot, x) = \int dx v_{t,x} \underbrace{[e^{ip \cdot \hat{x}}, N \gamma_t^{\text{HF}}](\cdot, x)}_{\text{gain } \hbar}.$$

Outline

- 1 Fermionic Many-Body Systems
- 2 Main Result: Hartree-Fock Theory is a Valid Approximation.
- 3 ... assuming certain Semi-Classical Estimates
- 4 Proof of Hartree-Fock Theory
- 5 Mixed States, Vlasov Equation**

Hartree-Fock for Mixed States

- e. g. positive temperature
- described by density matrices

$$\Gamma : \mathcal{F} \left(L^2(\mathbb{R}^3) \right) \rightarrow \mathcal{F} \left(L^2(\mathbb{R}^3) \right).$$

- **Purification**: density matrix \simeq vector in bigger Fock space

$$\mathcal{F} \left(L^2(\mathbb{R}^3) \oplus L^2(\mathbb{R}^3) \right).$$

- Araki-Wyss representation of the dynamics.
- Replace p-h transformation by general Bogoliubov transformation.

Vlasov Equation as Semi-Classical Limit

Hartree-Fock equation $\xrightarrow{\hbar \rightarrow 0}$ Vlasov equation

Summary

- Time-dependent Hartree-Fock theory is valid.
- Requires semi-classical structure (e. g. trapped ground states).
- Extension to mixed initial data; Derivation of Vlasov equation.