

VII LOCALIZATION OF BOUND STATES

We have seen: in asymptotically complete systems, the Hilbert space can be decomposed into a direct sum of bound states and scattering states, $\mathcal{H} = \mathcal{H}_B \oplus \text{ran } \Omega_+$, where $\mathcal{H}_B = \text{closure of } \{ \psi \in D(H) : H\psi = E\psi \}$. eigenvectors

Scattering states behave like free states: particle flies away and disperses.

$\because H = H^*, \sigma(H) \subset \mathbb{R}$, so $E \in \mathbb{R}$.

Eigenvectors (bound states) have trivial time evolution:
 $H\psi = E\psi \Rightarrow e^{-iHt}\psi = e^{-iEt}\psi$

The probability to find a particle near position $x \in \mathbb{R}^n$:
 $|e^{-iHt}\psi(x)|^2 = |\psi(x)|^2$ indep. of time.

Goal: Show that $|\psi(x)|^2 \simeq e^{-c|x|}$ for large $|x|$.

PROP.: (IMS localization formula)

Let $V: \mathbb{R}^n \rightarrow \mathbb{R}$ measurable and such that $H = -\frac{\Delta}{2} + V$ is closed with domain $H^2(\mathbb{R}^n)$.

Let $f \in C^\infty(\mathbb{R}^n)$, real-valued and $\partial^\alpha f \in L^\infty(\mathbb{R}^n) \forall |\alpha| \leq 2$.

Then on $H^2(\mathbb{R}^n)$ we have

$$2fHf = f^2H + Hf^2 + |\nabla f|^2$$

PROOF: We only show $2fHf\psi = (f^2H + Hf^2 + |0f|^2)\psi$ $\forall \psi \in C_0^\infty(\mathbb{R}^n)$.

The identity for $\psi \in H^2(\mathbb{R}^n)$ follows by an approximation argument in the graph norm of H .

(A priori it is not clear that $-\Delta = \mathcal{F}^{-1} p^2 \mathcal{F}$ satisfies a product rule; only when applied to C^2 -functions, where it agrees with the classical derivatives.)

Straightforward calculation:

$$2f(-\frac{\Delta}{2} + V)f\psi = f^2(-\frac{\Delta}{2} + V)\psi + (-\frac{\Delta}{2} + V)f^2\psi + |0f|^2\psi.$$

(because $f\psi \in D(V) \cap D(-\frac{\Delta}{2})$; and $fVf\psi$ can be written out pointwise almost everywhere.)

Use product rule for Laplacian:

$$\begin{aligned} \text{LHS: } f\Delta(f\psi) &= f\operatorname{div}(\operatorname{grad} f\psi + f\operatorname{grad}\psi) \\ &= f(\Delta f)\psi + f(\operatorname{grad} f)(\operatorname{grad}\psi) + f(\operatorname{div} f)(\operatorname{grad}\psi) + f^2(\Delta\psi). \end{aligned}$$

$$\begin{aligned} \text{RHS: } \frac{1}{2}f^2(\Delta\psi) + \frac{1}{2}\operatorname{div}(\operatorname{grad}(f^2\psi)) - |0f|^2\psi \\ = \dots = \text{LHS.} \end{aligned}$$

↙ Check it!



THM.: (Exp. localization of eigenvectors, Agmon)

Let $V \in L^2_{loc}(\mathbb{R}^n)$ and $n \geq 3$, or $V \in L^\infty(\mathbb{R}^n)$.

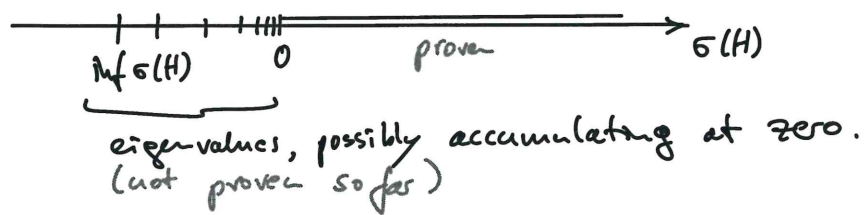
Let V real-valued and $V(x) \rightarrow 0$ ($|x| \rightarrow \infty$).

Let $H = -\frac{\Delta}{2} + V$.

If $\psi \in H^2(\mathbb{R}^n)$ and $H\psi = E\psi$, then

$e^{\beta|x|}\psi \in L^2(\mathbb{R}^n) \quad \forall \beta > 0$ with $E + \frac{\beta^2}{2} < 0$.

RMK: $\sigma(-\frac{\Delta}{2}) \subset \sigma(H)$ if Ω_\pm exist. $\sigma(-\frac{\Delta}{2}) = [0, \infty)$.



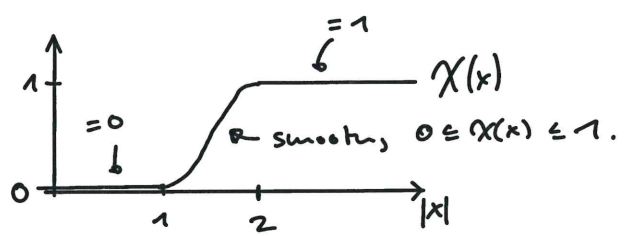
This means for the decay of eigenvectors, $H\psi = E\psi$:

$\psi(x) \sim e^{-\beta|x|}$ for $|x| \rightarrow \infty$,

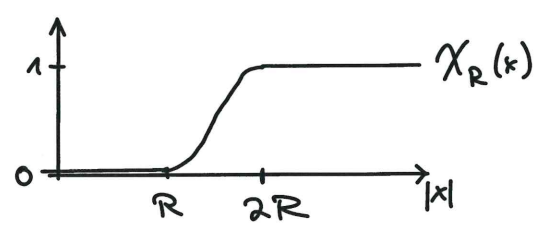
but rate of decay β smaller and smaller for E closer and closer to zero.

PROOF (of exp. localization):

Let $\chi \in C^\infty(\mathbb{R}^n)$:



Let $\chi_R(x) := \chi(\frac{x}{R})$:
(for $R > 0$)



Let $f(x) := \frac{\beta|x|}{1 + \epsilon|x|}$

(for $\epsilon > 0$ this is a regularization of $\beta|x|$).

Properties (check!):

$|f| \leq \frac{\beta}{2}$, and for $x \neq 0$: $|\nabla f(x)| \leq \beta$.
 (forgetting +1 in denominator) (growth \searrow slower than of $\beta|x|$, did the growth of $\beta|x|$ is proportional to β)

let $G := \chi_R e^f$. Then $G \in C^\infty(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$,

(*) $\nabla G = (\nabla \chi_R) e^f + \chi_R e^f (\nabla f) = (\nabla \chi_R) e^f + G (\nabla f) \in L^\infty(\mathbb{R}^n)$,
 and $\partial_i \partial_j G \in L^\infty(\mathbb{R}^n) \forall i, j \in \{1, \dots, n\}$.

Now: $\langle G \psi, (H-E) G \psi \rangle \stackrel{IMS}{=} \langle \psi, \left(\frac{1}{2} [G^2 H + H G^2 + |\nabla G|^2] - E \right) \psi \rangle$

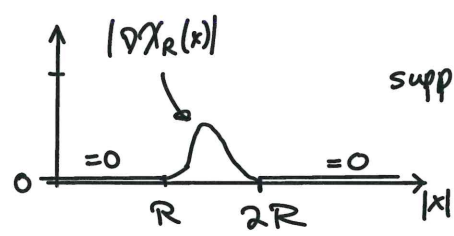
$H^* = H, H \psi = E \psi \Rightarrow \langle \psi, \frac{1}{2} |\nabla G|^2 \psi \rangle$

$\stackrel{(*)}{=} \frac{1}{2} \langle \psi, (|\nabla \chi_R|^2 e^{2f} + 2 \nabla \chi_R \cdot \nabla f e^f G + G^2 |\nabla f|^2) \psi \rangle$

← the other side.

$\Rightarrow \langle G \psi, (H-E - \frac{|\nabla f|^2}{2}) G \psi \rangle \leq \sup_{x \in \mathbb{R}^n} (|\nabla \chi_R|^2 e^{2f} + 2 |\nabla \chi_R| \beta e^{2f}) \|\psi\|_{L^2}^2$

Notice:



$\text{supp } |\nabla \chi_R(x)| \subset \{x \in \mathbb{R}^n : R \leq |x| \leq 2R\}$.

In $\text{supp } |\nabla \chi_R(x)|$: $e^{2\beta|x|} \leq e^{4\beta R} < \infty$.

$\Rightarrow \langle G \psi, (H-E - \frac{|\nabla f|^2}{2}) G \psi \rangle \leq C_R \|\psi\|_{L^2}^2, C_R < \infty. (i)$

On the other hand:

$$\langle \psi, (-\Delta)\psi \rangle = \int_{\mathbb{R}^n} |\hat{\psi}(p)|^2 p^2 \geq 0.$$

$$\langle G\psi, (H - E - \frac{10f^2}{2})G\psi \rangle \geq \langle G\psi, (V - E - \frac{10f^2}{2})G\psi \rangle$$

$$10f^2 \leq \delta \geq \langle G\psi, (V - E - \frac{\delta^2}{2})G\psi \rangle$$

$$\text{supp } \psi \subset \mathbb{R}^n \setminus B_R(0) \geq \left(\inf_{|x| \geq R} V(x) - E - \frac{\delta^2}{2} \right) \|G\psi\|_{L^2}^2$$

Recall: $E < 0$. > 0 for some $R > 0$ because $V(x) \rightarrow 0$ ($|x| \rightarrow \infty$).

$$\geq \delta_R \|G\psi\|_{L^2}^2. \quad (ii)$$

Combining (i) & (ii) $\Rightarrow \|G\psi\|_{L^2}^2 \leq \frac{C_R}{\delta_R} \|\psi\|_{L^2}^2$.

estimate uniform in $\varepsilon > 0!$

$$\Rightarrow \int_{|x| \geq 2R} e^{2\beta|x|} |\varphi(x)|^2 dx = \int_{|x| \geq 2R} \lim_{\varepsilon \downarrow 0} e^{2f(x)} |\varphi(x)|^2 dx$$

Lebesgue monotone convergence \rightarrow

$$= \lim_{\varepsilon \downarrow 0} \int_{|x| \geq 2R} e^{2f(x)} |\varphi(x)|^2 dx$$

$$\leq \lim_{\varepsilon \downarrow 0} \int \underbrace{N_R^2(x)}_{= G(x)^2} e^{2f(x)} |\varphi(x)|^2 dx$$

charact. fun. of the ball $B_{2R}(0)$ is bounded by $N_R(x)^2$ (look at the pictures above).

$$\leq \lim_{\varepsilon \downarrow 0} \frac{C_R}{\delta_R} \|\varphi\|^2 = \frac{C_R}{\delta_R} \|\varphi\|^2.$$

$$\int_{|x| \in 2R} e^{2\beta|x|} |\varphi(x)|^2 dx \leq \int_{|x| \in 2R} e^{4R\beta} |\varphi(x)|^2 dx$$

$$\leq e^{4R\beta} \|\varphi\|_{L^2}^2 < \infty.$$

$$\Rightarrow e^{\beta|x|} \varphi \in L^2(\mathbb{R}^n).$$



RMK: One can strengthen the result to $|\varphi(x)| \leq C_\beta e^{-\beta|x|}$ almost everywhere.

Lemma 7.3: Let $u \in L^2(\mathbb{R}^n)$ with $\hat{u} \in L^1(\mathbb{R}^n)$ and $\int |p| \hat{u}(p) dp < +\infty$. Then $u \in C^1(\mathbb{R}^n)$ and $\partial_{x_k} u(x) = (2\pi)^{-n/2} \int e^{-ip \cdot x} (-ip_k) \hat{u}(p) dp$.

Proof: Recall $u(x) = (2\pi)^{-n/2} \int e^{-ip \cdot x} \hat{u}(p) dp$ a.e.

We use the following standard corollary of Lebesgue's dominated convergence:

Let $f: \mathbb{R}^n \times \mathbb{R} \rightarrow \mathbb{C}$ such that $\partial_t f(x,t)$ exists and

(81)

there exists $g \in L^1(\mathbb{R}^n)$ such
that $|\partial_t f(x,t)| \leq g(x) \quad \forall x \in \mathbb{R}^n, t \in \mathbb{R}$.

Then $\int f(x,t) dx \rightarrow$ differentiable
w.r.t. t and

$$(*) \quad \frac{d}{dt} \int f(x,t) dx = \int \partial_t f(x,t) dx.$$

We take $f(p,x) := e^{-ip \cdot x} \hat{u}(p)$.

Then $\partial_x f(p,x) = e^{-ip \cdot x} (-ip) \hat{u}(p)$
exists and

$$|\partial_x f(p,x)| = |p| |\hat{u}(p)| =: g(p)$$

which \rightarrow by assumption an integrable
dominating function.

By (*) the partial derivatives exist
and are given by

$$\partial_k u(x) = (2\pi)^{-n/2} \int e^{-ip \cdot x} (-ip_k) \hat{u}(p) dp.$$

Continuity w.r.t. x follows also by dominated
convergence. ▀

Theorem 7.4: (Sobolev lemma)

(82)

Let $u \in H^s(\mathbb{R}^n)$, where $s > \frac{n}{2} + k$
for some $k \in \mathbb{N}$. Then:

(i) $u \in C^k(\mathbb{R}^n)$ and

$$\partial^\alpha u(x) \rightarrow 0 \quad (|x| \rightarrow \infty) \quad \text{for } |\alpha| \leq k.$$

$$(ii) \quad \sup_{|\alpha| \leq k} \|\partial^\alpha u\|_{L^\infty} \leq C_{s,k} \|u\|_{H^s}.$$

Proof: By induction, we generalize Lemma 7.3
to wave than one derivative: let $\alpha = (\alpha_1, \dots, \alpha_n)$

If $|p|^k \hat{u} \in L^1(\mathbb{R}^n)$ and $|\alpha| = \alpha_1 + \dots + \alpha_n \leq k$

$$\text{then } \partial^\alpha u(x) = (\partial_x)^{-|\alpha|} \int e^{-ip \cdot x} \underbrace{(\underbrace{i p}_1)^{\alpha_1} \dots (\underbrace{i p}_n)^{\alpha_n}}_{= (ip_1)^{\alpha_1} (ip_2)^{\alpha_2} \dots (ip_n)^{\alpha_n}} \hat{u}(p) dp.$$

Now let $u \in H^s(\mathbb{R}^n)$, with $s > \frac{n}{2} + k$.

Then $u \in L^2(\mathbb{R}^n)$ and

$$\begin{aligned} & \int |\hat{u}(p)| (1+p^2)^{k/2} dp \\ &= \int |\hat{u}(p)| (1+p^2)^{s/2} (1+p^2)^{\frac{k}{2} - \frac{s}{2}} dp \\ &\stackrel{CS}{\leq} \left(\int |\hat{u}(p)|^2 (1+p^2)^s dp \right)^{1/2} \left(\int \frac{1}{(1+p^2)^{s-k}} dp \right)^{1/2} \end{aligned}$$

$$= \|u\|_{H^s}$$

finite by
assumption

finite because
 $2(s-h) > n$

(go to spherical
coordinates)

83

So $|p|^k \hat{u} \in L^1(\mathbb{R}^n)$ as had to be
shown.

The Riemann-Lebesgue lemma
implies $\partial^\alpha u(x) \rightarrow 0$ ($|x| \rightarrow +\infty$).

This concludes (i).

By the same Cauchy-Schwarz as above:

$$\begin{aligned} |\partial^\alpha u(x)| &\leq (2\pi)^{-n/2} \int |e^{-i p \cdot x} (-i p)^\alpha \hat{u}(p)| dp \\ &\leq \|u\|_{H^s} (2\pi)^{-n/2} \left(\int \frac{1}{(1+p^2)^{s-h}} dp \right)^{1/2}. \end{aligned}$$

This concludes (ii).

